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## CHAPTER 6

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## 6.0 LAPLACE'S AND POISSON'S EQUATIONS AND UNIQUENESS THEOREM

- *In realistic electrostatic problems, one seldom knows the **charge distribution** – thus all the solution methods introduced up to this point have a limited use.*
- *These solution methods will not require the knowledge of the **distribution of charge**.*

# 6.1 LAPLACE'S AND POISSON'S EQUATIONS

To derive Laplace's and Poisson's equations, we start with **Gauss's law** in point form :

$$\nabla \cdot \bar{D} = \nabla \cdot \epsilon \bar{E} = \rho_v \quad (1)$$

Use **gradient** concept :

$$\bar{E} = -\nabla V \quad (2)$$

$$\begin{aligned} \nabla \cdot [\epsilon(-\nabla V)] &= \rho_v \\ \nabla \cdot \nabla V &= -\frac{\rho_v}{\epsilon} \end{aligned} \quad (3)$$

Operator :

$$\nabla \cdot \nabla = \nabla^2 \quad (4)$$

Hence :

$$\nabla^2 V = -\frac{\rho_v}{\epsilon} \quad V / m^2 \quad (5) \Rightarrow \text{Poisson's equation}$$

is called **Poisson's equation** applies to a **homogeneous media**.

**When the free charge density  $\rho_v = 0$**

$$\nabla^2 V = 0 \quad V / m^2 \quad (6) \quad \Rightarrow \text{Laplace's equation}$$

**In rectangular coordinate :**

$$\nabla^2 V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2}$$

## 6.2 UNIQUENESS THEOREM

**Uniqueness theorem states that for a  $V$  solution of a particular electrostatic problem to be unique, it must satisfy two criterion :**

- (i) Laplace's equation**
- (ii) Potential on the boundaries**

**Example :** In a problem containing two infinite and parallel conductors, one conductor in  $z = 0$  plane at  $V = 0$  Volt and the other in the  $z = d$  plane at  $V = V_0$  Volt, we will see later that the  $V$  field solution between the conductors is  $V = V_0 z / d$  Volt.

This solution will **satisfy Laplace's equation** and the **known boundary potentials** at  $z = 0$  and  $z = d$ .

Now, the  $V$  field solution  $V = V_0(z + 1) / d$  will **satisfy Laplace's equation** but will **not give the known boundary potentials** and thus is not a solution of our particular electrostatic problem.

Thus,  $V = V_0 z / d$  Volt is the only solution (**UNIQUE SOLUTION**) of our particular problem.

## 6.3 SOLUTION OF LAPLACE'S EQUATION IN ONE VARIABLE

**Ex.6.1:** Two infinite and parallel conducting planes are separated  $d$  meter, with one of the conductor in the  $z = 0$  plane at  $V = 0$  Volt and the other in the  $z = d$  plane at  $V = V_0$  Volt. Assume  $\rho_v = 0$  and  $\epsilon = 2\epsilon_0$  between the conductors.

Find : (a)  $V$  in the range  $0 < z < d$  ; (b)  $\bar{E}$  between the conductors ;  
(c)  $\bar{D}$  between the conductors ; (d)  $D_n$  on the conductors ; (e)  $\rho_s$  on the conductors ; (f) capacitance per square meter.

**Solution :**

(a) Since  $\rho_v = 0$  and the problem is in rectangular form, thus

$$\nabla^2 V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = 0 \quad (1)$$

**We note that  $V$  will be a function of  $z$  only  $V = V(z)$  ; thus :**

$$\nabla^2 V = \frac{\partial^2 V}{\partial z^2} = 0 \quad (2)$$

$$\frac{d^2 V}{dz^2} = \frac{d^2}{dz^2} \left( \frac{V}{dz} \right) = 0 \quad (3)$$

**Integrating twice :**

$$\frac{dV}{dz} = A \quad (4)$$

$$V = Az + B \quad (5)$$

**where  $A$  and  $B$  are constants and must be evaluated using given potential values at the boundaries :**

$$V|_{z=0} = B = 0 \quad (6)$$

$$\begin{aligned} V|_{z=d} &= Ad = V_0 \\ \rightarrow A &= V_0 / d \end{aligned} \quad (7)$$

**Substitute (6) and (7) into general equation (5) :**

$$\therefore V = \frac{V_0}{d} z \quad (V) \quad 0 < z < d$$

**(b)**

$$\begin{aligned} \bar{E} &= -\nabla V = -\left( \hat{x} \frac{\partial V}{\partial x} + \hat{y} \frac{\partial V}{\partial y} + \hat{z} \frac{\partial V}{\partial z} \right) \\ &= -\hat{z} \frac{\partial V}{\partial z} = -\hat{z} \frac{V_0}{d} \quad (V / m) \end{aligned}$$

**(c)**

$$\bar{D} = \epsilon \bar{E} = -\hat{z} \frac{2\epsilon_0 V_0}{d} \quad (C / m^2)$$



**(d) Surface charge :**

$$\begin{aligned}\rho_s|_{z=0} &= \bar{D} \cdot \hat{n} = -\hat{z} \frac{2\varepsilon_0 V_0}{d} \cdot \hat{z} \\ &= -\frac{2\varepsilon_0 V_0}{d}\end{aligned}$$

$$\begin{aligned}\rho_s|_{z=d} &= \bar{D} \cdot \hat{n} = -\hat{z} \frac{2\varepsilon_0 V_0}{d} \cdot (-\hat{z}) \\ &= +\frac{2\varepsilon_0 V_0}{d} \quad (C / m^2)\end{aligned}$$

**(e) Capacitance :**

$$\begin{aligned}C &= Q / V_{ab} \\ &= \frac{|\rho_s ds|}{V_0}\end{aligned}$$

$$\begin{aligned}\therefore C / m^2 &= \frac{|\rho_s|}{V_0} \\ &= \frac{2\varepsilon_0 V_0 / d}{V_0} = \frac{\varepsilon}{d} \quad (F / m^2)\end{aligned}$$



**Ex.6.2:** Two infinite length, concentric and conducting cylinders of radii  $a$  and  $b$  are located on the  $z$  axis. If the region between cylinders are charged free and  $\varepsilon = 3\varepsilon_0$ ,  $V = V_0$  (V) at  $a$ ,  $V = 0$  (V) at  $b$  and  $b > a$ . Find the capacitance per meter length.

**Solution :** Use Laplace's equation in cylindrical coordinate :

and  $V = f(r)$  only :

$$\nabla^2 V = \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial V}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 V}{\partial \phi^2} + \frac{\partial^2 V}{\partial z^2} = 0$$

*and  $V = f(r)$  only :*

$$\nabla^2 V = \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial V}{\partial r} \right) = 0$$

$$\frac{\partial}{\partial r} \left( r \frac{\partial V}{\partial r} \right) = 0$$

$$\left( r \frac{\partial V}{\partial r} \right) = A$$

$$\frac{\partial V}{\partial r} = \frac{A}{r}$$

$$V = A \ln r + B \quad (1)$$

**Boundary condition :**

$$V|_{r=a} = V_0 = A \ln a + B$$

$$V|_{r=b} = 0 = A \ln b + B$$

$$V = A \ln r + B \quad (1)$$

**Solving for A and B :**

$$A = \frac{V_0}{\ln(a/b)} \quad ; \quad B = \frac{-V_0 \ln b}{\ln(a/b)}$$

**Substitute A and B in (1) :**

$$\therefore V = \frac{V_0 \ln(b/r)}{\ln(b/a)} \quad ; \quad a < r < b$$

$$\bar{E} = -\nabla V = -\hat{r} \frac{\partial V}{\partial r} = \frac{V_0}{r \ln(b/a)} \hat{r}$$

$$\bar{D} = \epsilon \bar{E} = \frac{\epsilon V_0}{r \ln(b/a)} \hat{r}$$

$$\therefore V = \frac{V_0 \ln(b/r)}{\ln(b/a)}$$

**Surface charge densities:**

$$\rho_s|_{r=a} = \bar{D} \cdot \hat{r} = \frac{\epsilon V_0}{a \ln(b/a)}$$

$$\rho_s|_{r=b} = \bar{D} \cdot -\hat{r} = -\frac{\epsilon V_0}{b \ln(b/a)}$$

**Line charge densities :**

$$\rho_\ell|_{r=a} = \rho_s|_{r=a} (2\pi a) = \frac{2\pi\epsilon V_0}{\ln(b/a)}$$

$$\rho_\ell|_{r=b} = \rho_s|_{r=b} (2\pi b) = -\frac{2\pi\epsilon V_0}{\ln(b/a)}$$

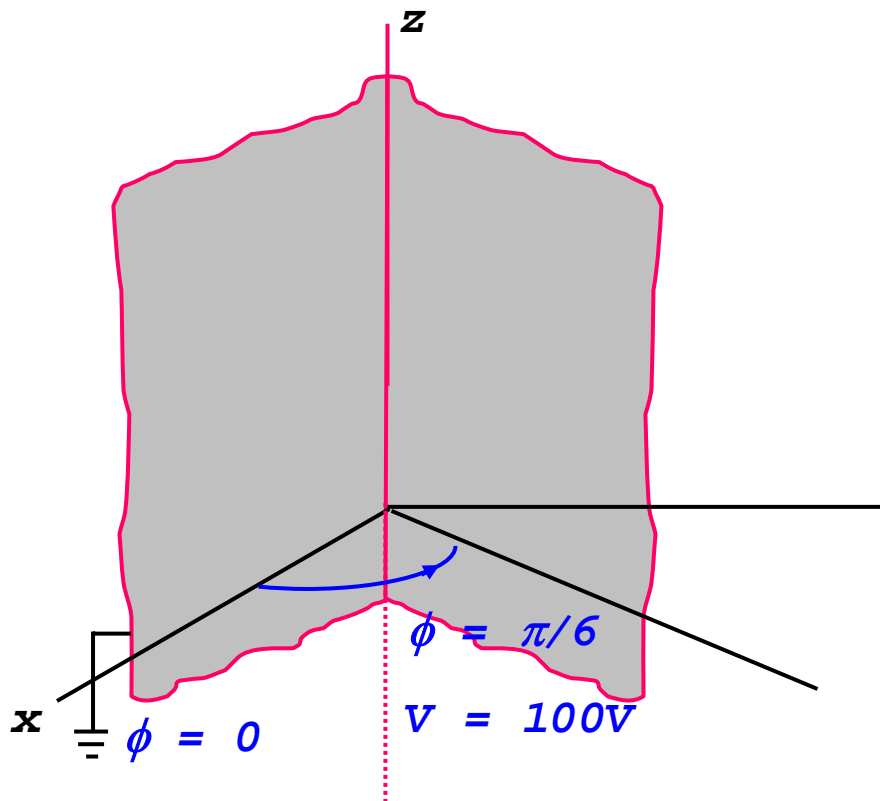
**Capacitance per unit length:**

$$C = \frac{Q}{V_{ab}} = \frac{|\rho_\ell d\ell|}{V_o}$$

$$C/m = \frac{|\rho_\ell|}{V_o} = \frac{2\pi\epsilon}{\ln(b/a)} \quad (F/m)$$

**Ex.6.3:** Two infinite conductors form a wedge located at

$\phi = 0$  and  $\phi = \pi / 6$  is as shown in the figure below. If this region is characterized by charged free. Find  $V$  and  $\vec{E}$ . Assume  $V = 0$  V at  $\phi = 0$  and  $V = 100$  V at  $\phi = \pi / 6$ .



**Solution :**  $V = f(\phi)$  in cylindrical coordinate :

$$\nabla^2 V = \frac{1}{r^2} \frac{d^2 V}{d\phi^2} = 0$$

$$\frac{d^2 V}{d\phi^2} = 0$$

$$\frac{dV}{d\phi} = A$$

$$V = A\phi + B$$

**Boundary condition :**

$$V|_{\phi=0} = 0 = B$$

$$V|_{\phi=\pi/6} = 100 = A(\pi/6)$$

$$A = 600/\pi$$

**Hence :**

$$V = \frac{600}{\pi} \phi$$

$$\begin{aligned} \bar{E} &= -\nabla V = -\frac{1}{r} \frac{dV}{d\phi} \hat{\phi} \\ &= -\frac{600}{\pi r} \hat{\phi} \end{aligned}$$

**for region :**  
 $0 \leq \phi \leq \pi/6$



**Ex.6.4:** Two infinite concentric conducting cone located at  $\theta = \pi/10$  and  $\theta = \pi/6$ . The potential  $V = 0$  V at  $\theta = \pi/10$  and  $V = 50$  V at  $\theta = \pi/6$ . Find  $V$  and  $\vec{E}$  between the two conductors.

**Solution :**  $V = f(\theta)$  in spherical coordinate :

$$\nabla^2 V = \frac{1}{r^2 \sin \theta} \frac{d}{d\theta} \left[ \sin \theta \frac{dV}{d\theta} \right] = 0$$

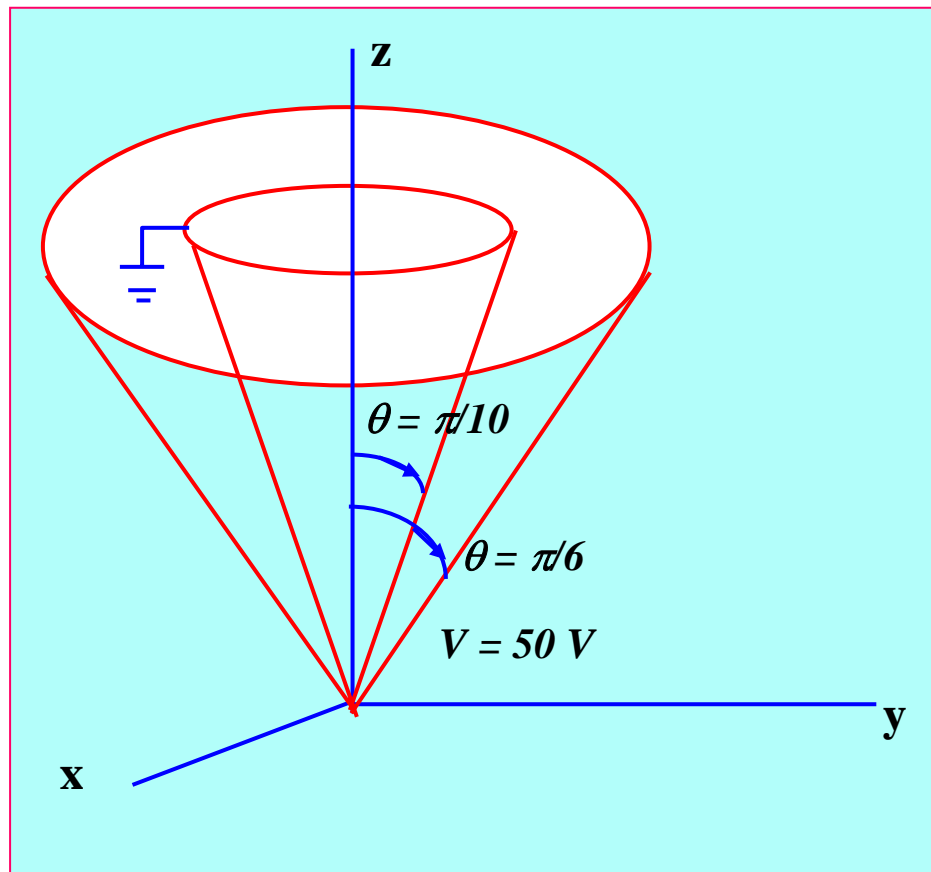
$$\frac{d}{d\theta} \left[ \sin \theta \frac{dV}{d\theta} \right] = 0$$

$$\sin \theta \frac{dV}{d\theta} = A$$

$$\frac{dV}{d\theta} = \frac{A}{\sin \theta}$$

$$V = A \ln(\tan \theta / 2) + B$$

Using :  $\int \frac{d\theta}{\sin \theta} = \ln(\tan \theta / 2)$



**Boundary condition :**

$$V = A \ln(\tan \theta / 2) + B$$

$$V \Big|_{\theta=\pi/10} = 0 = A \ln(\tan \pi / 20) + B$$

$$V \Big|_{\theta=\pi/6} = 50 = A \ln(\tan \pi / 12) + B$$

**Solving for A and B :**

$$A = \frac{50}{\ln\left(\frac{\tan \pi / 12}{\tan \pi / 20}\right)} ; B = -\frac{50 \ln \tan \pi / 20}{\ln\left(\frac{\tan \pi / 12}{\tan \pi / 20}\right)}$$

**Hence at region :**  $\theta/10 \leq \theta \leq \theta/6$

$$\begin{aligned} V &= \frac{50}{\ln\left(\frac{\tan \pi / 12}{\tan \pi / 20}\right)} \ln \frac{\tan \theta / 2}{\tan \pi / 20} \\ &= 95.1 \ln\left(\frac{\tan \theta / 2}{0.1584}\right) \end{aligned}$$

**and**

$$\begin{aligned} \bar{E} &= -\nabla V = \frac{1}{r} \frac{dV}{d\theta} \hat{\theta} \\ &= -\frac{95.1}{r \sin \theta} \hat{\theta} \end{aligned}$$

## 6. 4 SOLUTION FOR POISSON'S EQUATION

When the free charge density  $\rho_v \neq 0$

**Ex.6.5:** Two infinite and parallel conducting planes are separated  $d$  meter, with one of the conductor in the  $x = 0$  plane at  $V = 0$  Volt and the other in the  $x = d$  plane at  $V = V_0$  Volt. Assume  $\rho_v \neq 0$  and  $\varepsilon = 4\varepsilon_0$  between the conductors.

Find : (a)  $V$  in the range  $0 < x < d$  ; (b)  $\bar{E}$  between the conductors

**Solution :**

$V = f(x) :$

$$\nabla^2 V = -\frac{\rho_v}{\varepsilon}$$

$$\frac{d^2 V}{dx^2} = -\frac{\rho_0}{\varepsilon}$$

$$\frac{dV}{dx} = -\frac{\rho_0}{\varepsilon} x + A$$

$$V = -\frac{\rho_0}{\varepsilon} \frac{x^2}{2} + Ax + B$$

**Boundary condition :**

$$V|_{x=0} = 0 = B$$

$$V|_{x=d} = V_0 = -\frac{\rho_0}{\epsilon} \frac{d^2}{2} + Ad$$

$$A = \frac{V_0}{d} + \frac{\rho_0}{\epsilon} \frac{d}{2}$$

$$V = -\frac{\rho_0}{\epsilon} \frac{x^2}{2} + Ax + B$$

**In region :  $0 \leq x \leq d$**

$$V = \frac{\rho_0}{\epsilon} \frac{x}{2} (d - x) + \frac{V_0}{d} x$$

$$\bar{E} = -\frac{dV}{dx} \hat{x}$$

$$= -\left[ \frac{V_0}{d} + \frac{\rho_0}{\epsilon} \left( \frac{d}{2} - x \right) \right] \hat{x}$$

**Ex.6.6:** Repeat Ex.6.5 with  $\rho_v = 0$  and  $\epsilon_r = 1 + x$

**Solution :**

$$\nabla \cdot \overline{D} = \rho_v$$

$$\nabla \cdot \epsilon \overline{E} = 0$$

$$\frac{d}{dx} (1 + x) \epsilon_0 \overline{E} = 0$$

$$\frac{d}{dx} (1 + x) (-\nabla V) = 0$$

$$(1 + x) \left( -\frac{dV}{dx} \right) = A$$

$$\left( -\frac{dV}{dx} \right) = \frac{A}{1 + x}$$

$$V = -A \ln(1 + x) + B$$

**Boundary condition :**

$$V|_{x=0} = 0 = B$$

$$V|_{x=d} = V_0 = -A \ln(1+d)$$

$$\rightarrow A = -\frac{V_0}{\ln(1+d)}$$

**In region :**  $0 \leq x \leq d$

$$V = -A \ln(1+x) + B$$

$$V = V_0 \frac{\ln(1+x)}{\ln(1+d)}$$

$$\bar{E} = -\frac{dV}{dx} \hat{x} = -\frac{V_0}{(1+x)\ln(1+d)} \hat{x}$$